

A Conflict Between Chip Architecture and a Direction of Ocean Model Development

Alistair Adcroft
Princeton University

Atmosphere, Oceans, and Computational Infrastructure

California Institute of Technology

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When will my models become non-competitive?

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- General vertical coordinate algorithms, and new [physical] parameterizations of sub-grid scale processes, look more and more like column-wise code and some are non-vectorizable
 - This is non-optimal on existing architectures but worse moving forward
- GPUs offer significant speedups only on parts of the code
 - Effective speedup of the whole model is moderate, if even positive
- The general coordinate algorithms allow comparably very large time-steps which [currently] far out pace the benefits of GPUs
 - Are there ways to get the new algorithms to work on GPUs or TPUs?



Time-step and efficiency limitations

- Sound waves are filtered out
- Surface gravity waves
 - $c_{bt} \sim \sqrt{gH} \sim 200$ m/s
- Internal waves
 - $c_{iw} \sim NH \sim 2$ m/s
- Currents
 - $U \sim 1-3$ m/s
 - $W_{eddy} \sim 10$ m/day
 - $W_{IG\ waves} \sim 100$ m/day
 - $W_{overtURNS} \sim U \sim 1$ cm/s

	100 km	25 km	1 km	100 m
$\Delta x / c_{bt}$	8 minutes	2 minutes	5 seconds	½ second
$\Delta x / c_{iw}$	12 hours	3 hours	8 minutes	4 minutes
$\Delta x / U$	1 day	2 hours	5 minutes	2 minutes
f^{-1}	1.9 hours	1.9 hours	1.9 hours	1.9 hours
	50 m	10 m	1 m	1 cm
$\Delta z / W$	1 days	2 hours	15 minutes	10 seconds

- At coarse resolutions, numerics are dominated by considerations of rotational dynamics
- At fine resolutions, codes/methods should look more towards CFD
- **External mode continues to be a critical barrier**



Barotropic solver

- Equations are not up for debate!
 - Need to be solved with particular methods for consistency
 - **Discrepancies give tsunamis instead of tides**
- Small amount of computation
 - Almost linear
- Frequent communication
- Latency sensitive

$$\left. \begin{aligned} \partial_t u_k &= -\frac{1}{\rho_0} \nabla p + \dots \\ \partial_t h_k &= \nabla \cdot (h_k u_k) \end{aligned} \right\} \begin{array}{l} \text{Solve with baroclinic} \\ \Delta t_{bc} \end{array}$$

SSH (constraint):

$$\partial_t \eta = \sum_k \partial_t h_k = \nabla \cdot \sum_k h_k u_k = \nabla \cdot (HU)$$

Reconcile with:

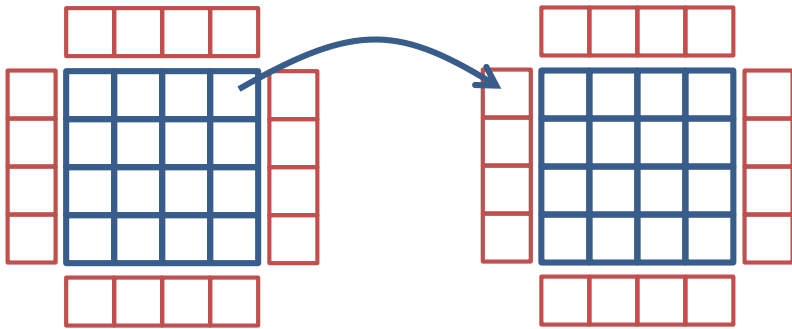
$$\left. \begin{aligned} \partial_t \eta &= \nabla \cdot (HU) \\ \partial_t U &= -g \nabla \eta + \dots \end{aligned} \right\} \begin{array}{l} \text{Solve with barotropic} \\ \Delta t_{bt} \end{array}$$

$$\Delta t_{bt} \ll \Delta t_{bc}$$

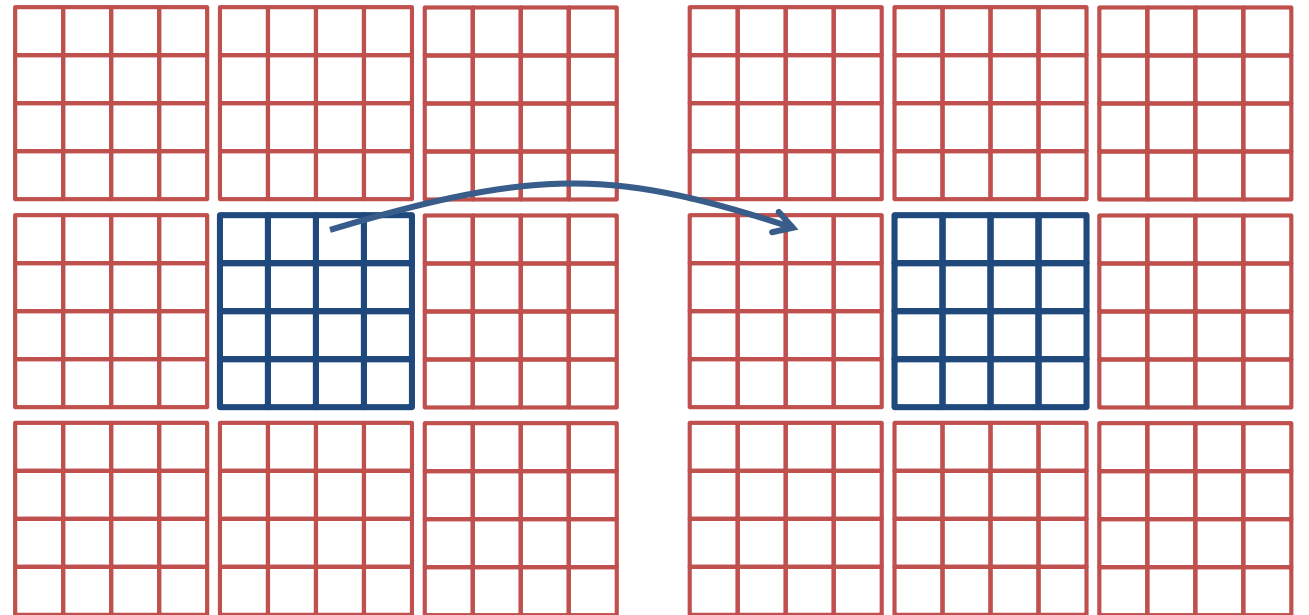


Adjustable strategies for barotropic solver

- Low latency/low bandwidth
 - Send small packets frequently



- High bandwidth/high latency
 - Send large packets infrequently

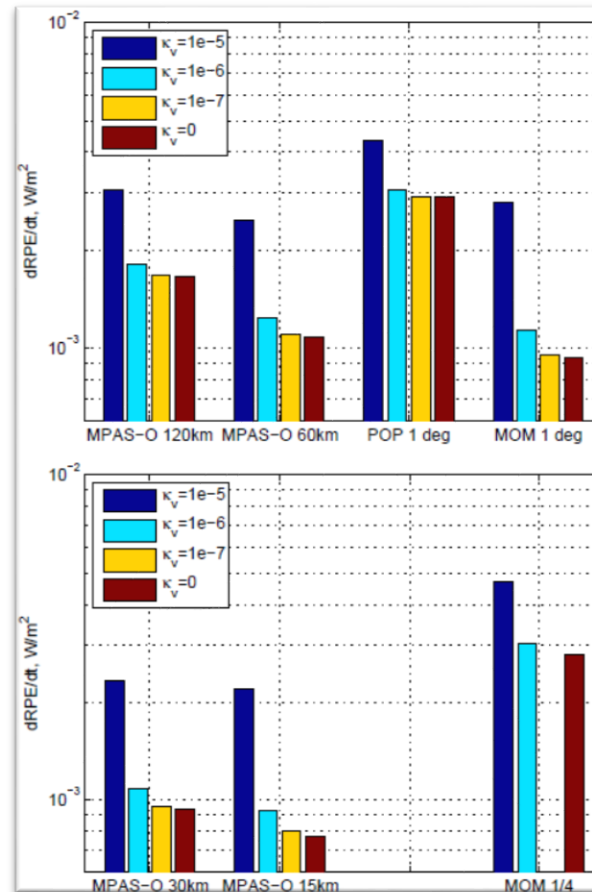
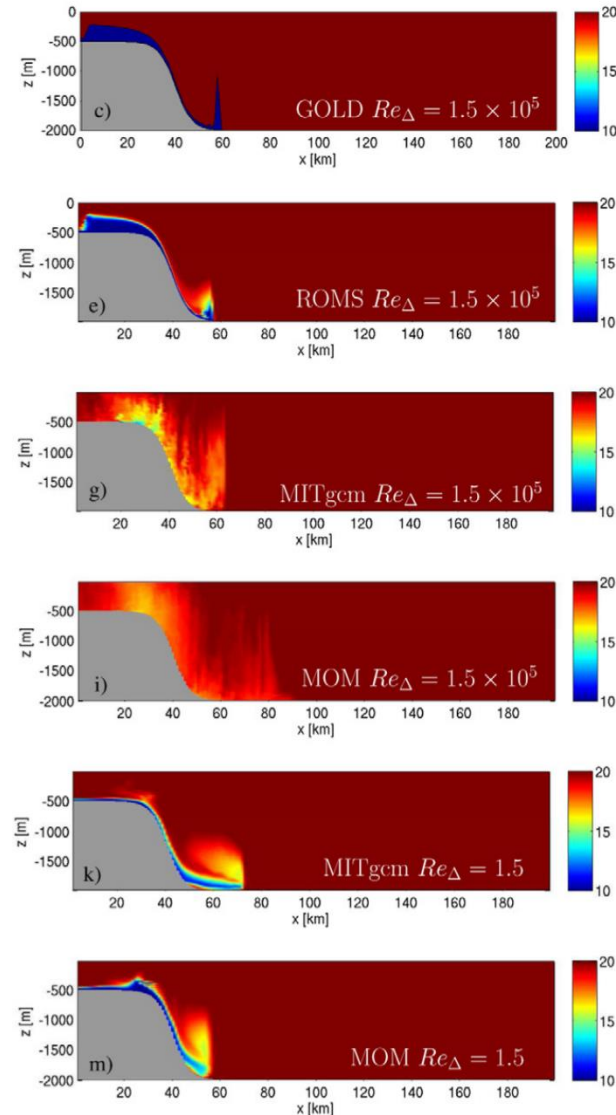


- Signal propagates across many cores in a single baroclinic time step

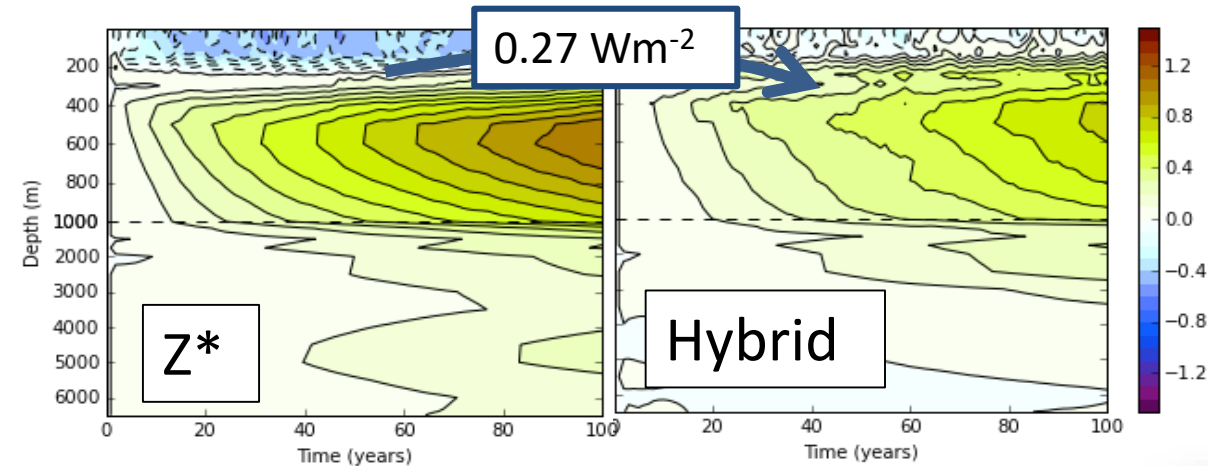
Vertical coordinates and spurious mixing

Coordinate MIP

Ilicak et al., 2012;
Petersen et al., 2014

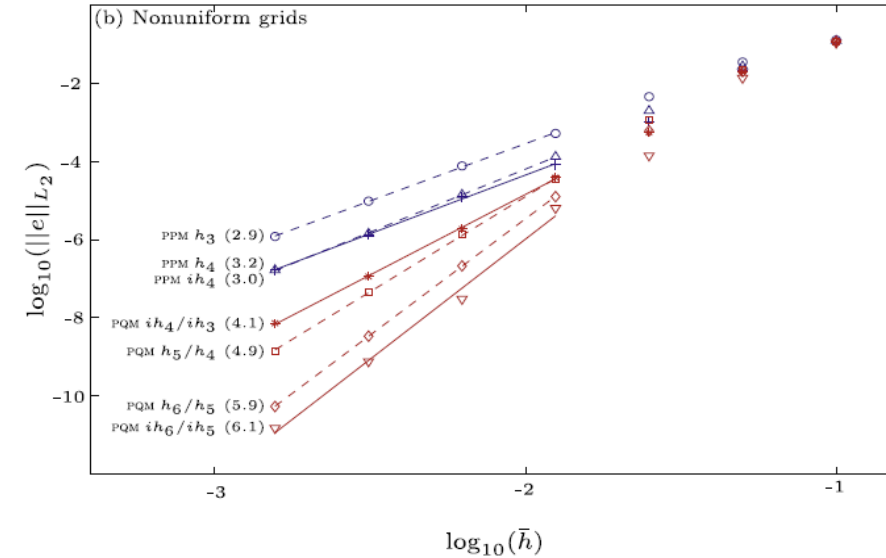
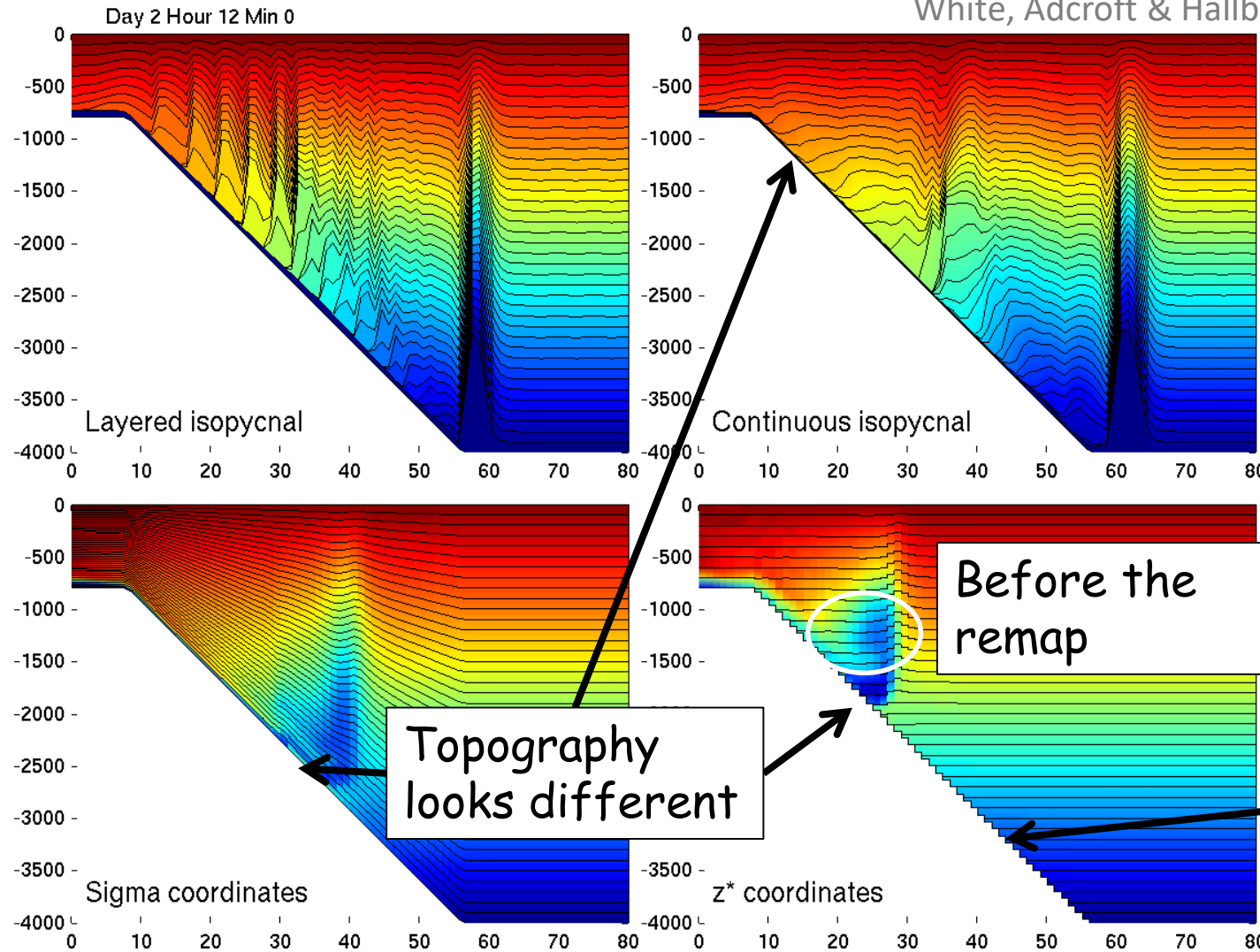


- Long debate about appropriate vertical coordinate
 - Now fairly well understood
- Spurious mixing:
 - grid Reynolds number + ...
 - long time scale questions



A.L.E. allows you to work with any coordinate

White, Adcroft & Hallberg, JCP 2009



Eulerian

$$\partial_z p = -g\rho(z, S^n, \theta^n)$$

$$v_h^{n+1} = v_h^n + \Delta t \left(-\frac{1}{\rho_o} \nabla_z p + \dots \right)$$

$$\partial_z w = -\nabla \cdot v_h^{n+1}$$

$$\theta^{n+1} = \theta^n - \Delta t \left[\begin{array}{l} \nabla \cdot (v_h^{n+1} \theta^n) + \\ \partial_z (w \theta^n) + \dots \end{array} \right]$$

$$\boxed{\frac{\Delta t w}{\Delta z} < 1}$$

A.L.E. (flavor 1)

$$\partial_z p = -g\rho(z, S^n, \theta^n)$$

$$v_h^{n+1} = v_h^n + \Delta t \left(-\frac{1}{\rho_o} \nabla_z p + \dots \right)$$

$$\delta_k (w^* + w_g) = -\nabla \cdot h^n v_h^{n+1}$$

$$h^{n+1} = h^n + \Delta t \delta_k (w_g)$$

$$h^{n+1} \theta^{n+1} = h^n \theta^n$$

$$-\Delta t \left[\begin{array}{l} \nabla \cdot (h^n v_h^{n+1} \theta^n) \\ + \delta_k (w^* \theta^n) + \dots \end{array} \right]$$

$$\boxed{\frac{\Delta t w^*}{\Delta z} < 1}$$

$$w^* = w - w_g$$

Leclair & Madec, 2011, use this form

A.L.E. (flavor 2)

$$\partial_z p = -g\rho(z, S^n, \theta^n)$$

$$v_h^\dagger = v_h^n + \Delta t \left(-\frac{1}{\rho_o} \nabla_z p + \dots \right)$$

$$h^\dagger = h^n - \Delta t \nabla \cdot (h^n v_h^\dagger)$$

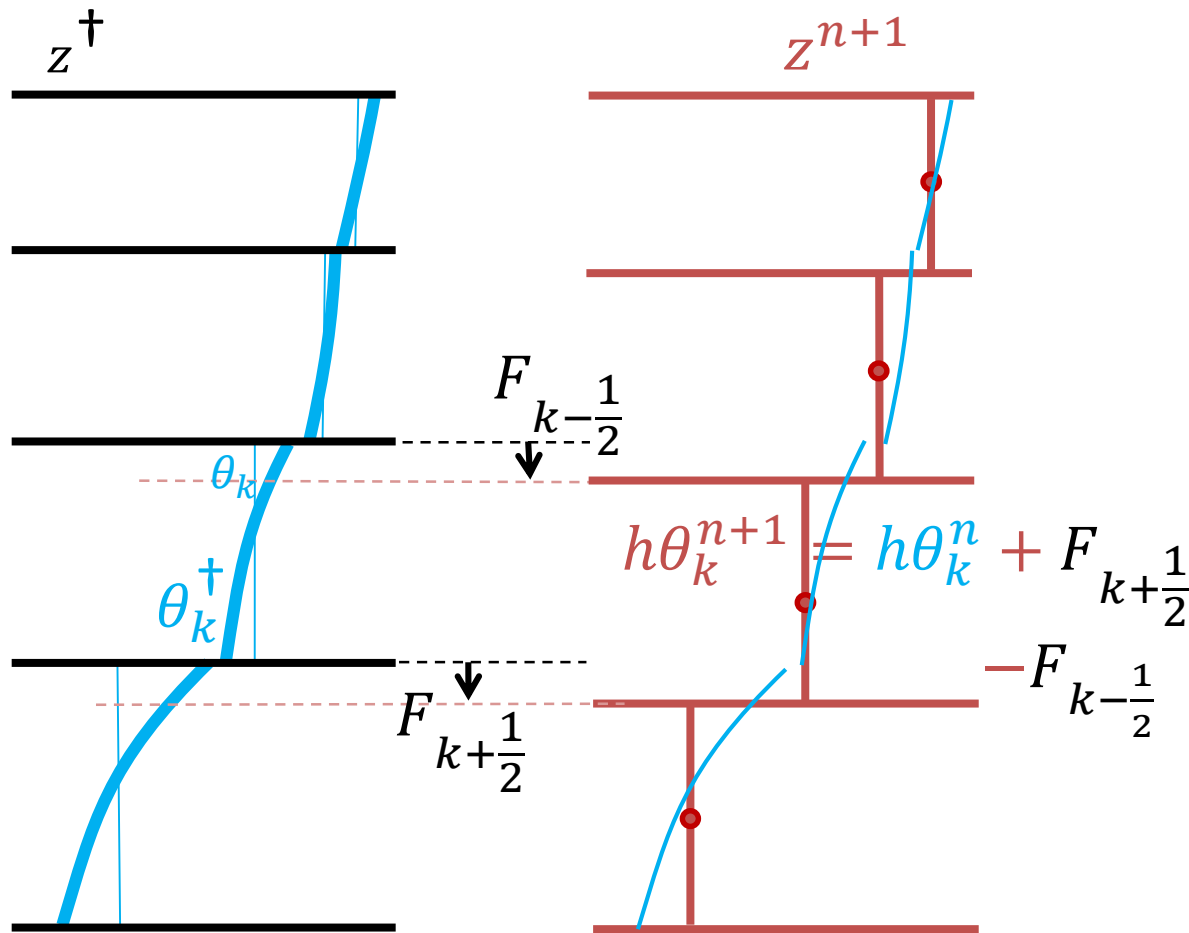
$$h^{n+1} \leftarrow \delta_k Z(z^\dagger)$$

$$\theta^{n+1} = \theta^\dagger (Z(z^\dagger))$$

Bleck, 2002



Remapping styles: flux form small CFL



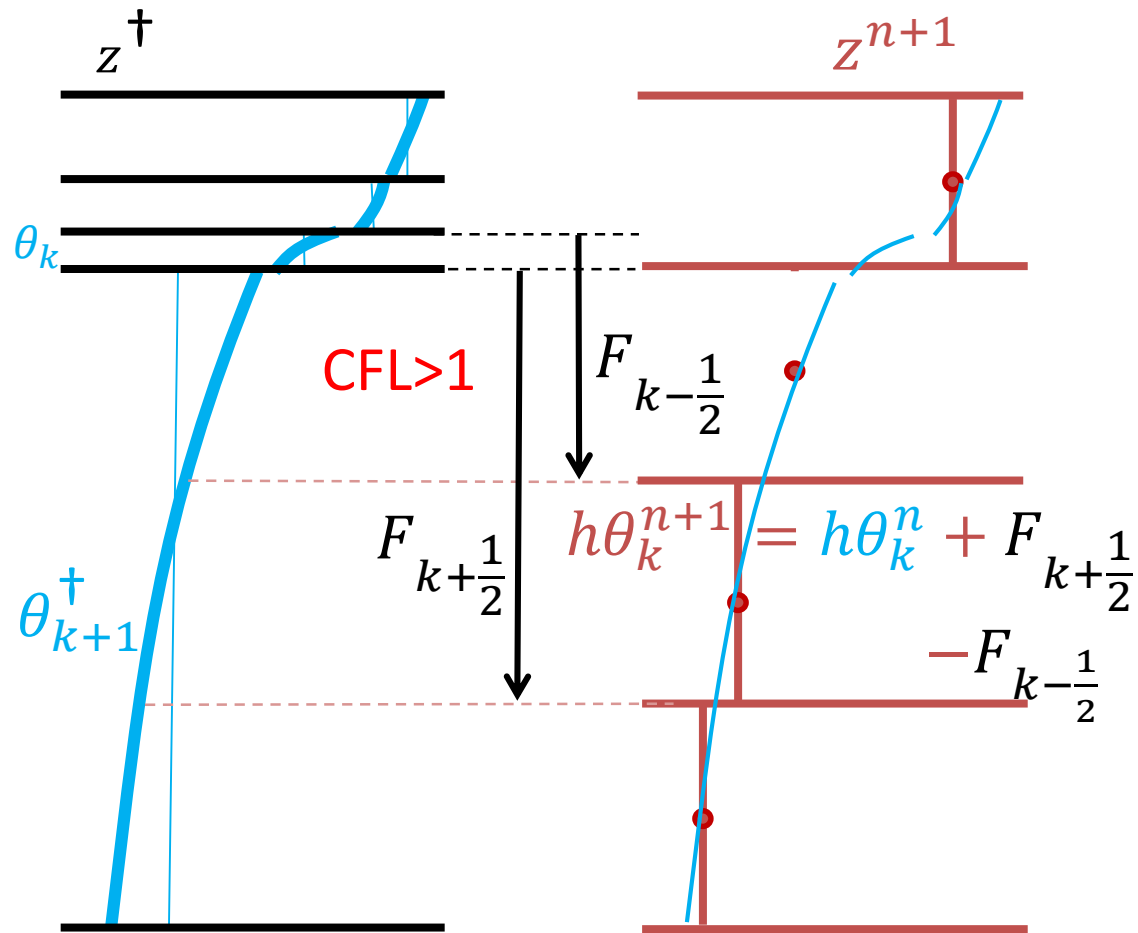
- Flux-form \rightarrow conservative
- CFL < 1 local stencil
- Stencil is “known”
e.g. The kernel

$$F_{k+\frac{1}{2}} = \max\left(0, \Delta t w_{k-\frac{1}{2}}\right) \theta_{k+1} + \min\left(0, \Delta t w_{k-\frac{1}{2}}\right) \theta_k$$

is vectorizable in horizontal

$$F[:, K] = \max(0, w[:, K]) \theta[:, k + 1] + \min(0, w[:, K]) \theta[:, k]$$

Remapping styles: flux form for arbitrary CFL

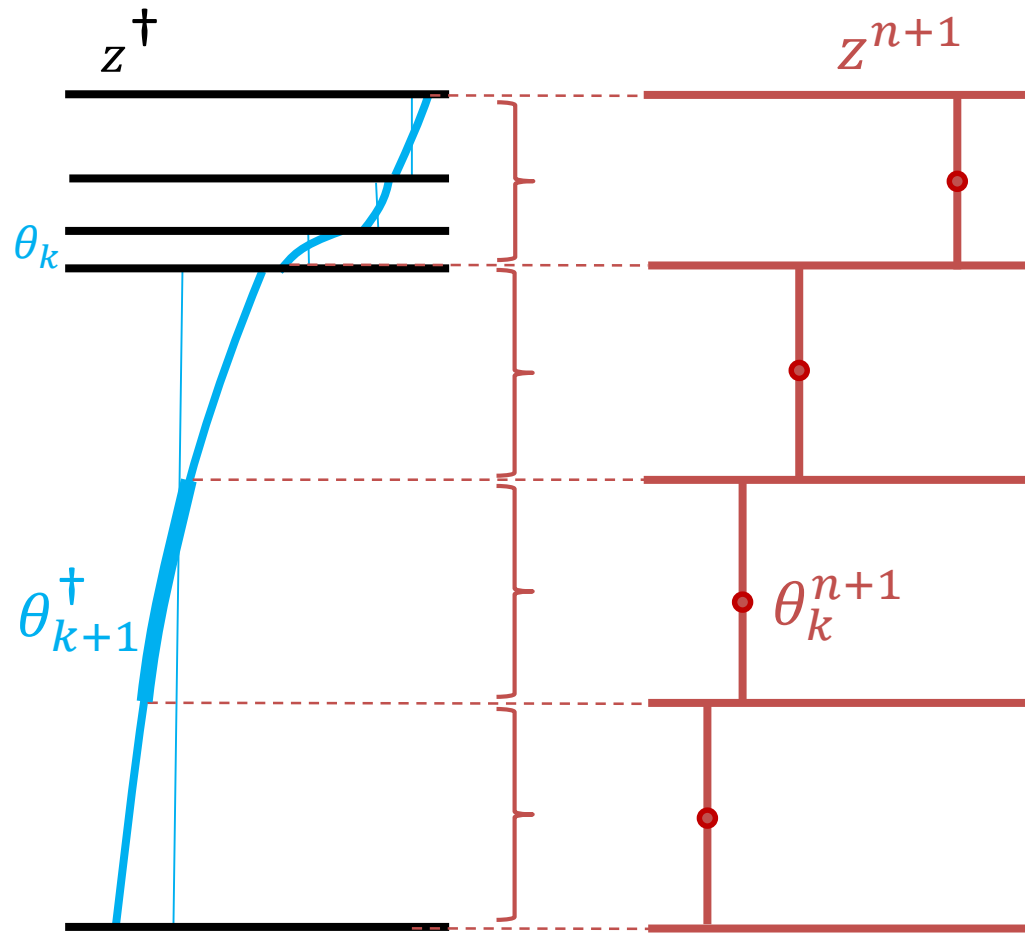


- Flux-form \rightarrow conservative
- CFL > 1 often evaluates **residual difference of large numbers**
 - Potentially inaccurate
- Stencil can be full column:

$$F_{k+1/2} = \begin{cases} \sum_k^{m-1} h_l \theta_l + \Delta t w_{k-1/2} \theta_m & w_{k-1/2} > 0 \\ \sum_{m+1}^{k-1} h_l \theta_l + \Delta t w_{k-1/2} \theta_m & w_{k-1/2} < 0 \end{cases}$$

- Nested loops – eek!

Remapping styles: “projection”

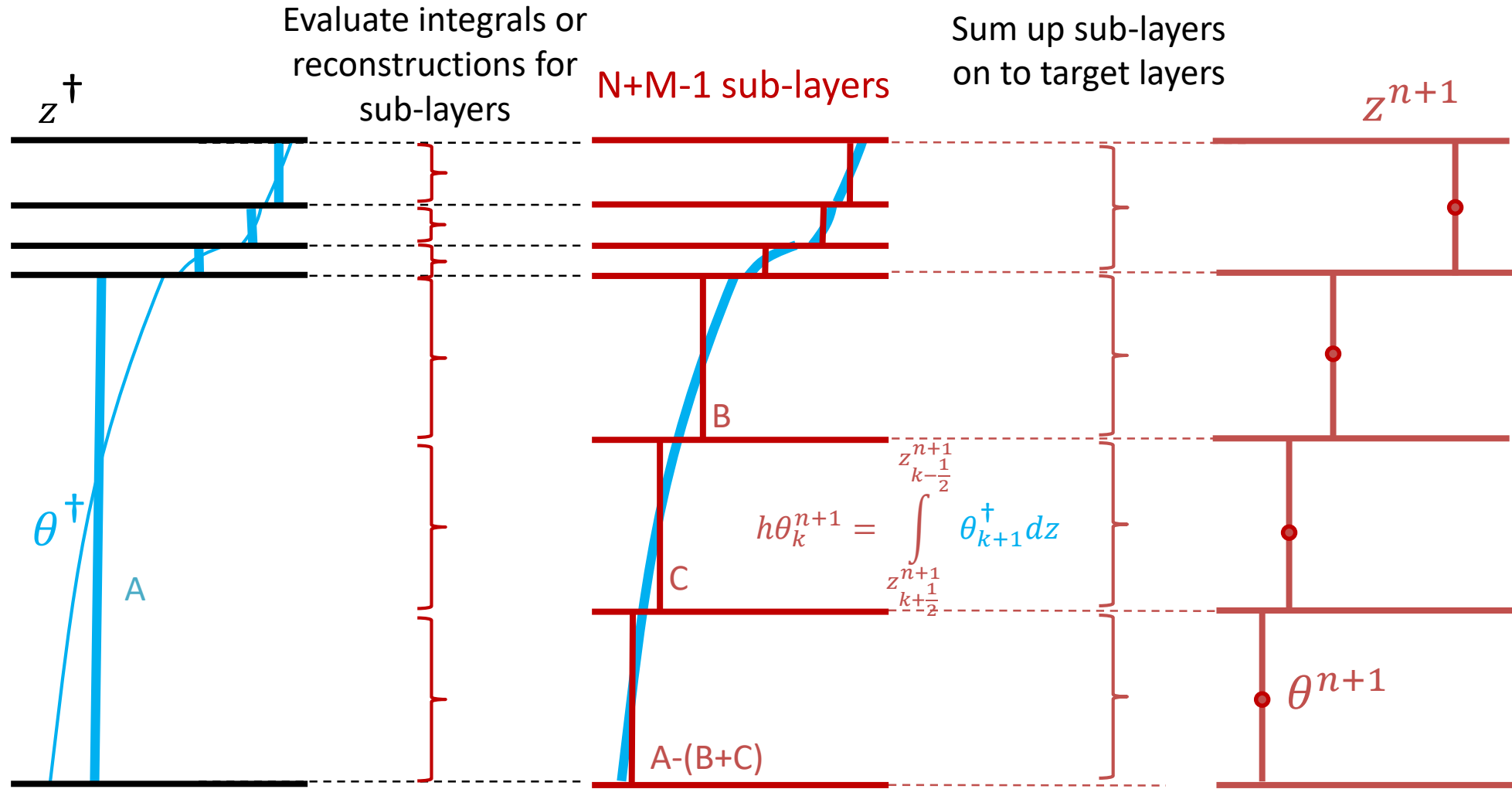


- Accuracy independent of CFL
 - No residual of differences
- Stencil can still be full column

$$h\theta_k^{n+1} = \int_{z_{k+\frac{1}{2}}^{n+1}}^{z_{k-\frac{1}{2}}^{n+1}} \theta_{k+1}^\dagger dz$$

- Conservation is less accurate
 - No equal/opposite flux terms

How we implement remapping in MOM6

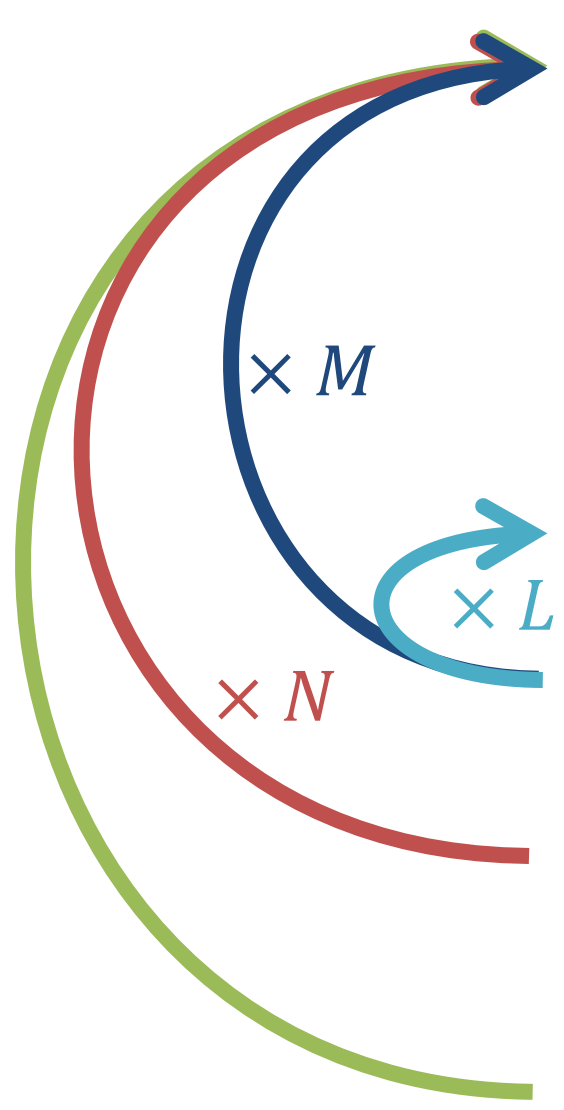


- No nested loops but there are now extra branches

Hallberg and Adcroft, in prep.



Subcycling to minimize compute per step



$$\delta_k p = -\rho(z, S^n, \theta^n) \delta_k \Phi$$

$$v_h^{m+1} = v_h^m + \frac{1 \Delta t}{M \rho_0} (-\nabla_r p - \rho \nabla_r \Phi + \dots)$$

$$h^{m+1} = h^m - \frac{1}{M} \Delta t \nabla_r \cdot (h^m v_h^{m+1})$$

Internal gravity waves

$$\frac{\Delta t c_g}{\Delta x} < 1$$

$$U^{l+1} = U^l + \frac{1}{L} \Delta t (-\nabla \eta^l + \dots)$$

$$\eta^{l+1} = \eta^m - \frac{1}{L} \Delta t \nabla_r \cdot (H U^{l+1})$$

Barotropic gravity waves

$$h^* \theta^* = h^n \theta^n - M \Delta t \left[\nabla \cdot \left(\sum_{m=1}^M h^m v_h^{m+1} \theta^n \right) \right]$$

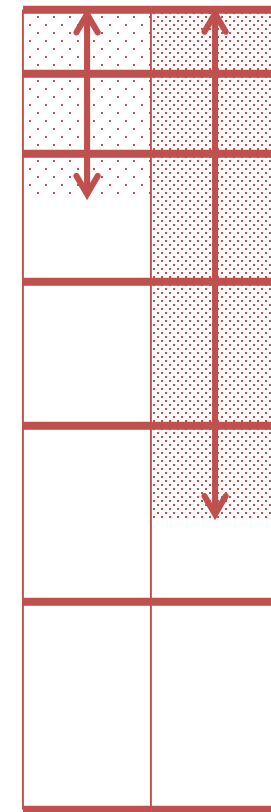
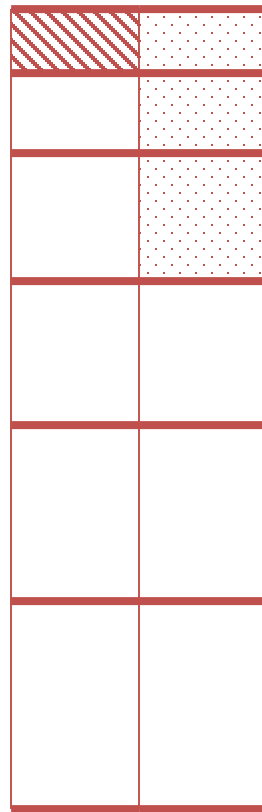
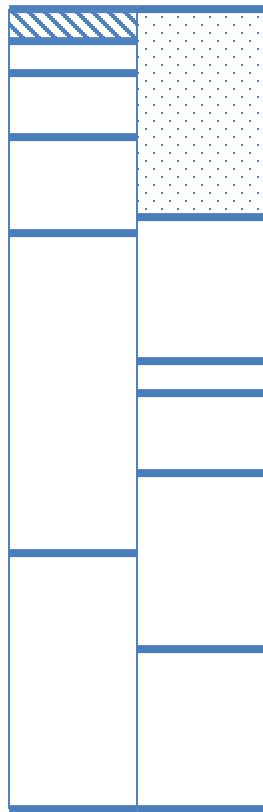
$$\frac{M \Delta t u_h}{\Delta x} < 1$$

$$h^{n+1} \leftarrow \delta_k Z(z^*); \theta^{n+1} = \theta^*(Z(z^*)); \dots$$



Non SIMD aspect of GC + column physics

- Column-to-column variations in source/target grid
- Some parameterizations have built-in non-linear non-local solvers



What about TPUs (ML)?

- All discretizations ultimately look like a [sparse] matrix multiply
 - Challenge for our algorithm is construction of the matrix in each column is very “branchy”
 - Different pattern of addressing in each column
 - Changes every step
- Could we delegate to ML?
 - I have not tried this but willing to entertain
 - However, ML is inexact

- Sea-level calculation:

$$\eta = -D + \sum_k h_k$$

$$\text{err}(\eta) \sim \epsilon D$$

- $\epsilon_{16\text{-bit}} = 2^{-11}$, $\text{err}(\eta)_{16\text{-bit}} \sim 3.4 \text{ m}$
- $\epsilon_{32\text{-bit}} = 2^{-24}$, $\text{err}(\eta)_{32\text{-bit}} \sim 4 \times 10^{-4} \text{ m}$ (0.4 mm)
- $\epsilon_{64\text{-bit}} = 2^{-53}$, $\text{err}(\eta)_{64\text{-bit}} \sim 8 \times 10^{-13} \text{ m}$ (.008 Å)
- Context of evolving sea-level:
 - Sea-level rise $\sim .3 \text{ m/Cy}$ or $1 \times 10^{-5} \text{ m/day}$
 - Tides $\sim 10 \text{ m / day}$



- Algorithm developments have yielded great gains
 - Challenging to implement on existing SIMD tech
- New hardware still doesn't address old problems
 - Barotropic solve likely to always be a performance barrier unless model fits on one chip
- Inexact math of ML is a reality
 - Potential advantages where stochasticity is appropriate
 - Major problem for algorithms that rely on consistency